

# Today - Lecture 13

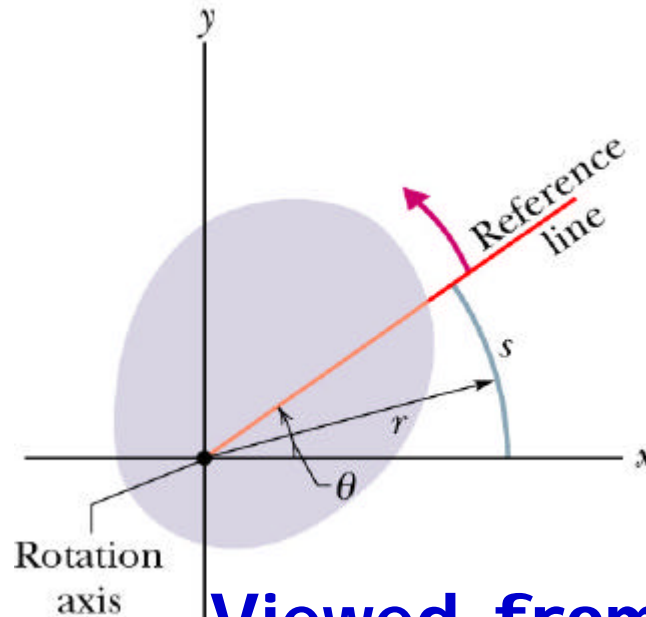
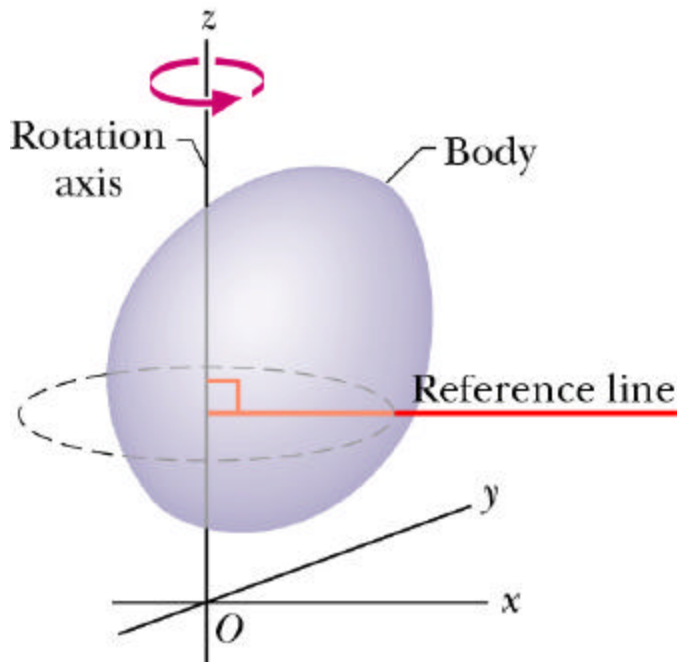
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- Today's lecture ... continue with rotations, torque, ...
- Note that chapters 11, 12, 13 all involve rotations

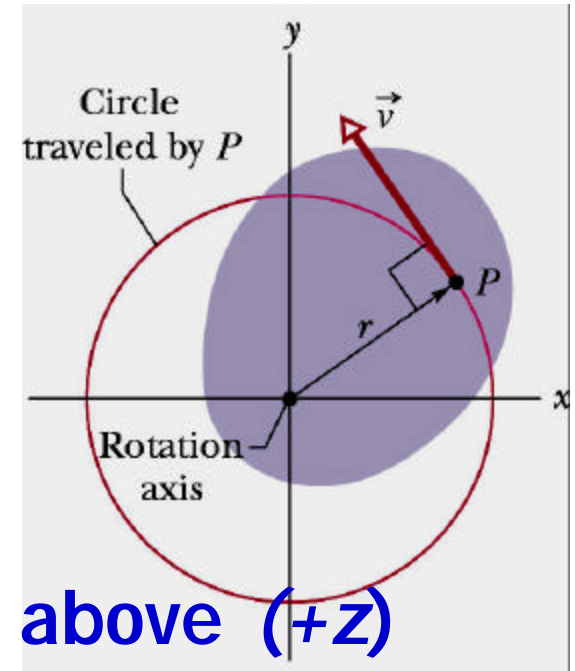
review

# Rotations

## Chapters 11 & 12



Viewed from above (+z)

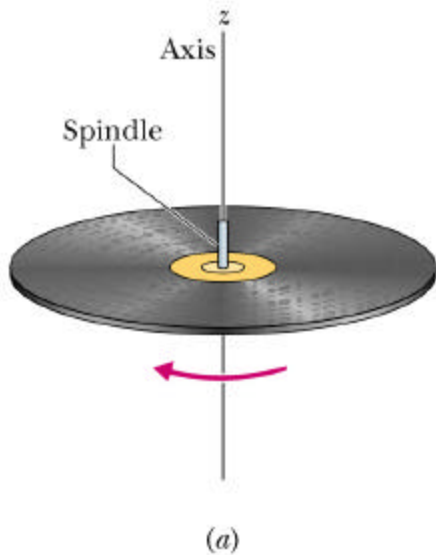


- Rotational, or angular velocity, gives tangential velocity

$$w \equiv \frac{dq}{dt} \quad v = \frac{ds}{dt} = r \frac{dq}{dt}$$

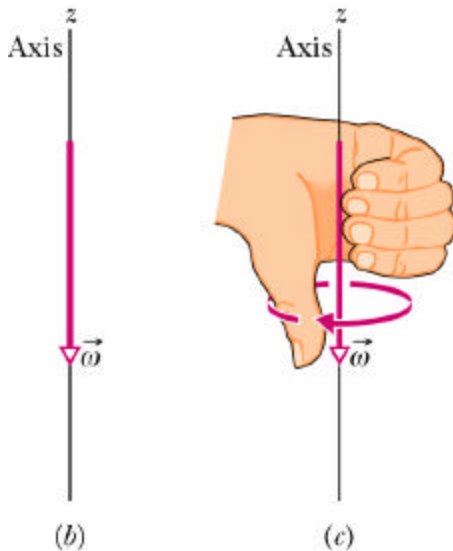
$$v = rw$$

# Vector angular velocity



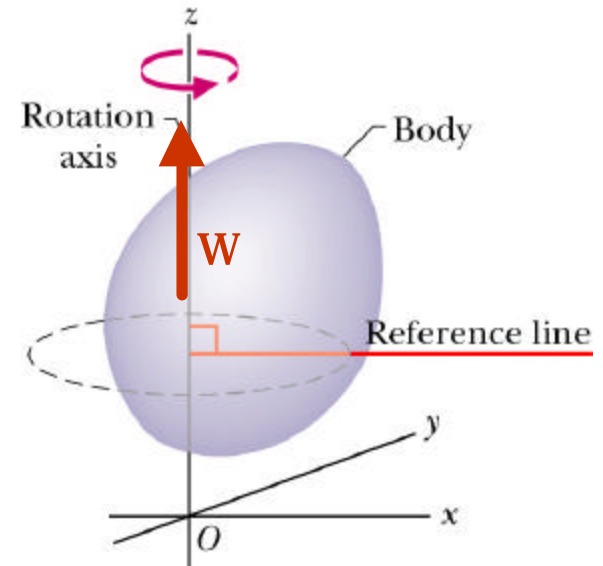
Define vector ang. velocity

$$\vec{\omega} \equiv \left\{ \begin{array}{l} \text{magnitude } \omega = \frac{dq}{dt} \\ \text{direction along rotation axis} \\ \text{sense using right-hand rule} \end{array} \right\}$$

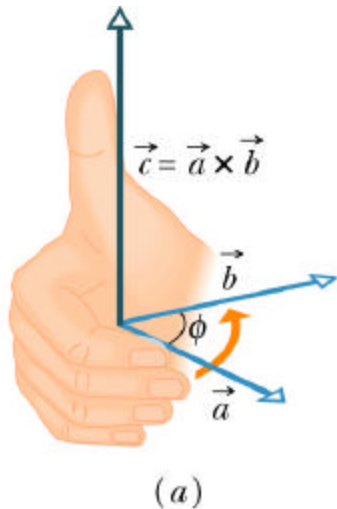


## Examples

- ◆ Record at left has  $\omega$  along  $-z$  axis
- ◆ body to right has  $\omega$  along  $+z$  axis

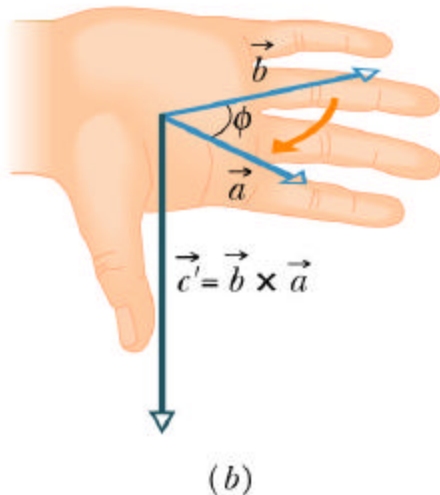


# Vector Product (ch 3)



Convention for  $\omega$  recalls the definition of the vector product (not a coincidence)

The vector product of chapter 3 is a convenient way to use the vector angular velocity (along with other rotation quantities)

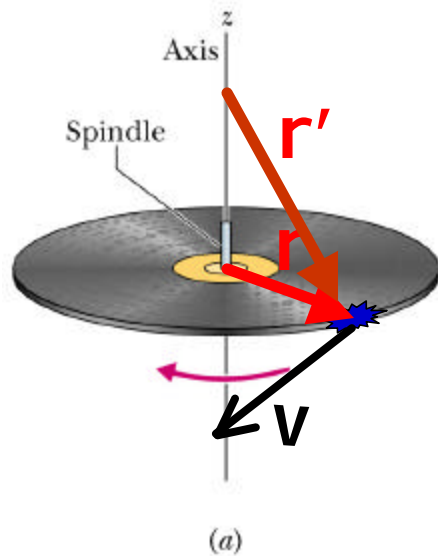


$$\vec{c} = \vec{a} \times \vec{b}$$

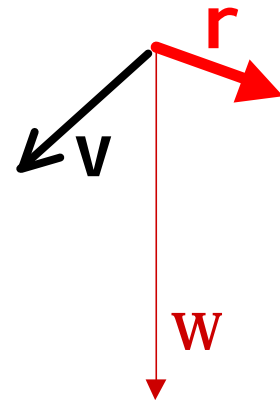
$$= -\vec{b} \times \vec{a}$$

$$c = ab \sin f$$

# Vector product gives tangential velocity



$$\mathbf{v} = r\omega$$

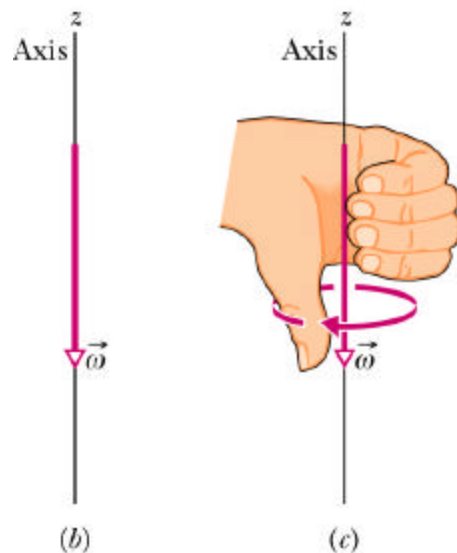


$$\vec{\mathbf{v}} = \vec{\mathbf{w}} \times \vec{\mathbf{r}}$$

Note rel part of  $r$  is the  $\wedge$  distance to the axis of rotation

$$\vec{\mathbf{v}} = \vec{\mathbf{w}} \times \vec{\mathbf{r}}$$

- Already seen the tangential speed in terms of dist from axis and ang velocity
- Vector or cross-product gives direction of  $v \wedge w \& r$
- True for any rotating vector



# One Step Further

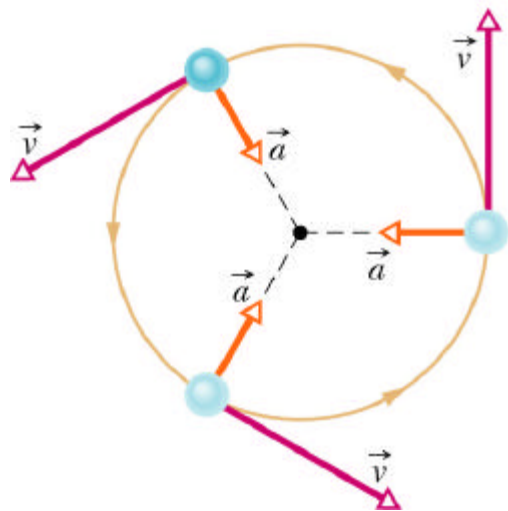
$$\vec{v} = \vec{\omega} \times \vec{r}$$

• Any rotating vector of fixed length simple time rate of change

- ◆ Was true for radius vector
- ◆ Also true for velocity vector

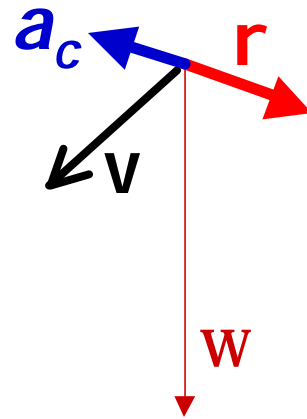
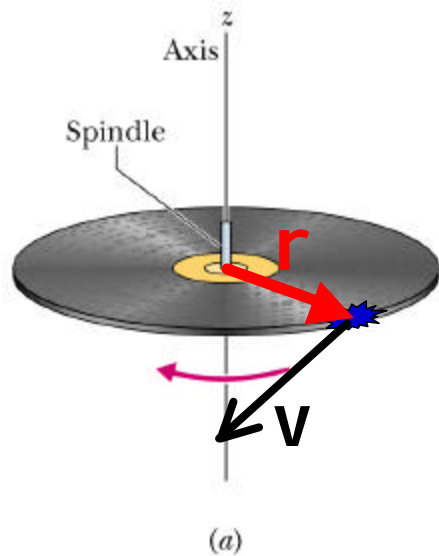
For  $\vec{c}$  any vector of const. length rotating  $\vec{\omega}$ , then

$$\frac{d\vec{c}}{dt} = \vec{\omega} \times \vec{c}$$



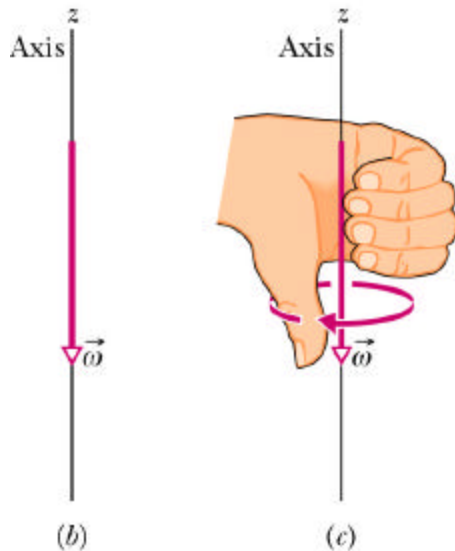
- $a_c = v^2/r$
- Directed inward
- The velocity vector
  - ◆ Has fixed magnitude
  - ◆ Rotates with angular velocity,  $\omega$

# Centripetal Acceleration (revisited as a Vector)



$$\vec{v} = \vec{\omega} \times \vec{r}$$

- Rate of change of position around axis given by cross-product
- Also, rate of change of velocity around axis of rotation given by same cross-product



$$\vec{a}_c = \vec{\omega} \times \vec{v}$$

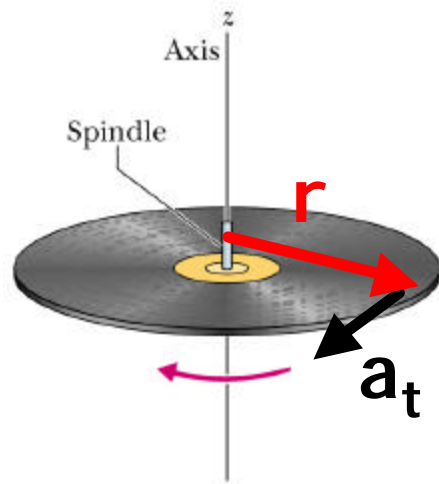
$$= \vec{\omega} \times (\vec{\omega} \times \vec{r})$$

$$\vec{a}_c = -\omega^2 \vec{r}$$

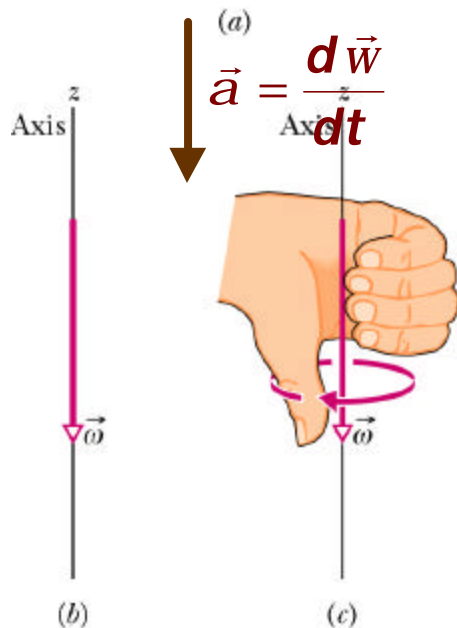
Vector Identity (appx E)

$$\vec{a} \times (\vec{b} \times \vec{c}) = \vec{b}(\vec{a} \cdot \vec{c}) - \vec{c}(\vec{a} \cdot \vec{b})$$

# Angular acceleration



- If ang velocity,  $\omega$ , changes, there is an angular acceleration,  $\alpha$ .
- Point at radius,  $r$ , has linear acceleration,  $a_t = r\alpha$ , tangent to circle at  $r$ .



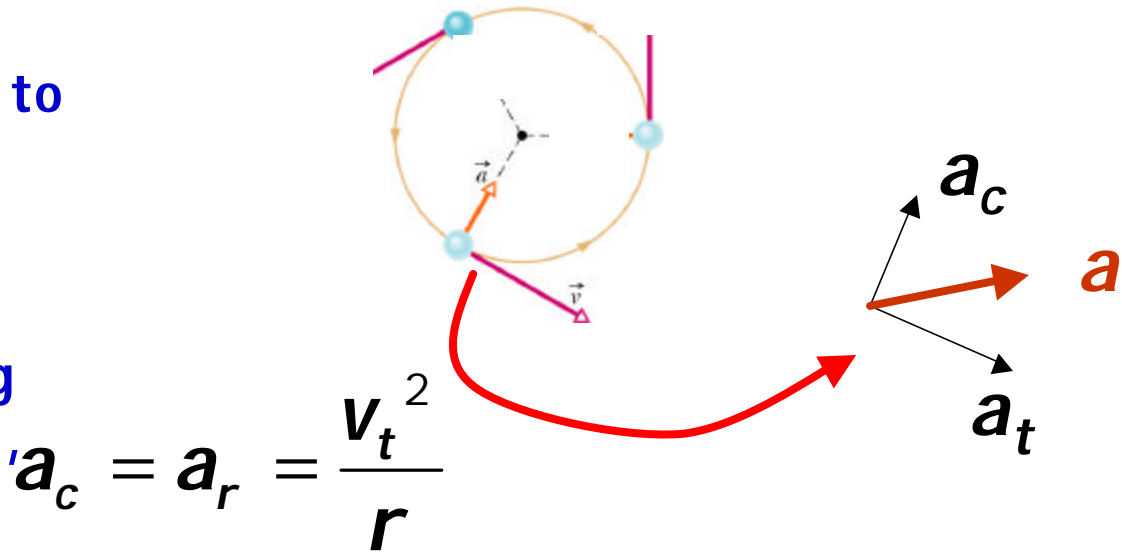
$$\vec{\alpha} \equiv \frac{d\vec{\omega}}{dt} \quad \vec{a}_t = \vec{\alpha} \times \vec{r}$$



angular acceleration.MOV

# Angular and Centripetal Acceleration

- Relationships for constant angular acceleration similar to those of constant linear acceleration
- See table 11-1
- Note that pt moving on circle, in general, will have two components of acceleration
- The two add vectorially to give total acceleration vector



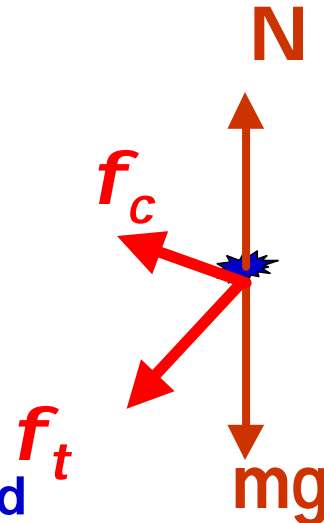
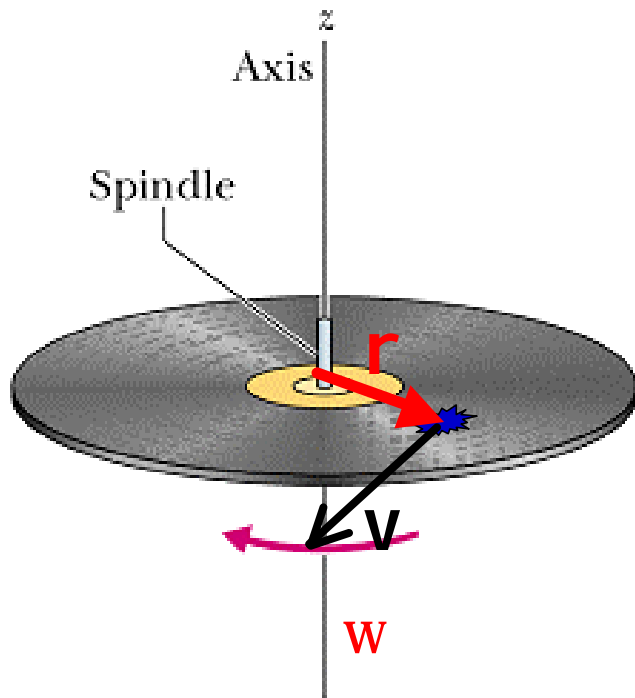
$$a_c = a_r = \frac{v_t^2}{r}$$

pointing to rotation axis

$$a_t = \frac{dv_t}{dt} = r a$$

along direction of rotation

# Hint: HW problem 11-32



- Watermelon seed on record player
  - ◆ Starts at rest  $w_0=0$
  - ◆ Reaches 33 1/3 rpm
- Convert to standard units
- Two components (at right angles) of acceleration
- Find net acceleration ... net friction force causes it

$$v = rw$$

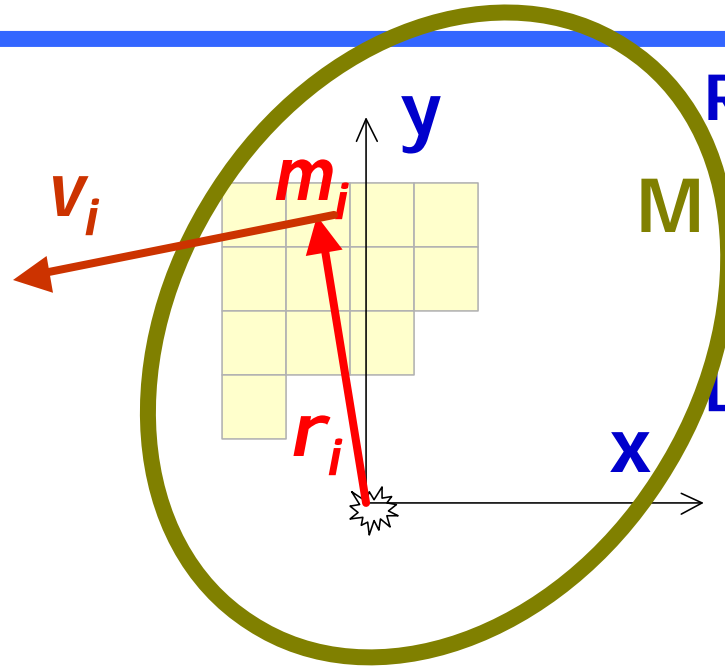
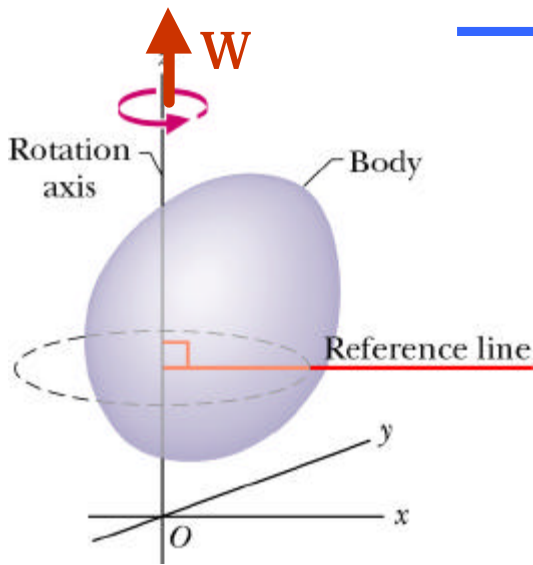
$$|a_c| = \frac{v^2}{r} = rw^2$$

$$|a_t| = ra$$

$$a = \frac{w - w_0}{\Delta t}$$

- Coefficient of static friction must be big enough to withstand this total force

# Kinetic Energy of Rotation: Moment of Inertia



Rotation of a rigid body about an axis implies Kinetic Energy...  
Look down from +z axis at mass pieces in body

$$K = \sum_i \frac{1}{2} m_i v_i^2$$

$$= \frac{1}{2} \sum_i m_i r_i^2 \omega^2$$

$$K = \frac{1}{2} I \omega^2$$

$$I \equiv \sum_i m_i r_i^2 \quad \text{discrete } m_i$$

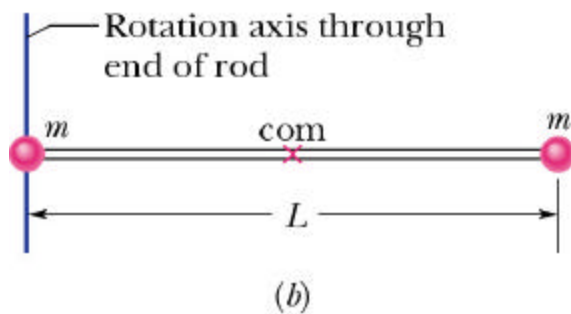
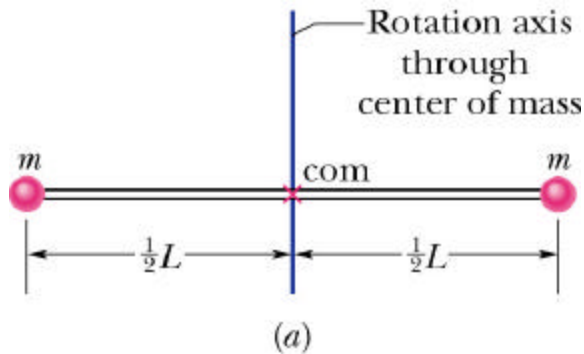
$$I \equiv \int_M r^2 dm \quad \text{continuous } m$$



energy in rotational motion.MOV

# Point masses: sample prob. 11-5

- Rigid body of two masses ( $m$ ) connected by rigid massless bar.
- Find  $I$  for two axes



$$I \equiv \sum_i m_i r_i^2 \quad \text{discrete } m_i$$

$$I \equiv \int_M r^2 dm \quad \text{continuous } m$$

$$I_{com} = m\left(\frac{1}{2}L\right)^2 + m\left(\frac{1}{2}L\right)^2$$

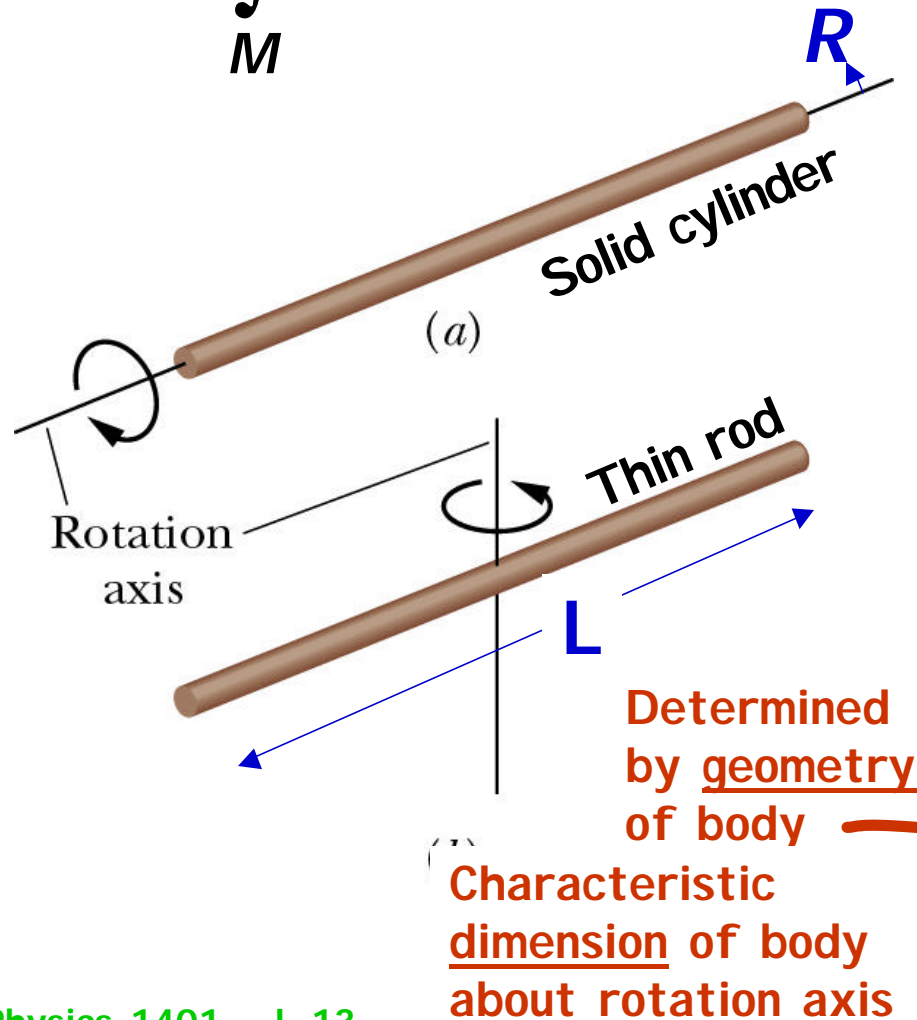
$$I_{com} = \frac{1}{2}mL^2$$

$$I_{end} = m(0)^2 + m(L)^2$$

$$I_{end} = mL^2$$

# Moments of Inertia ... continuous rigid bodies

$$I \equiv \int_M r^2 dm$$



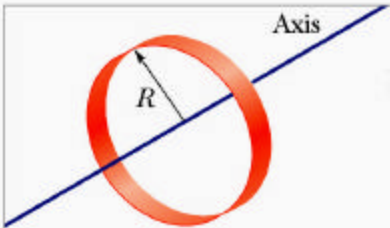
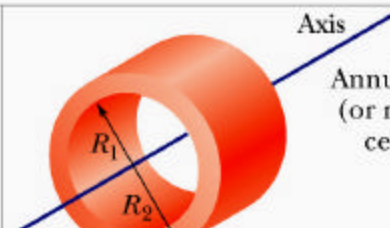
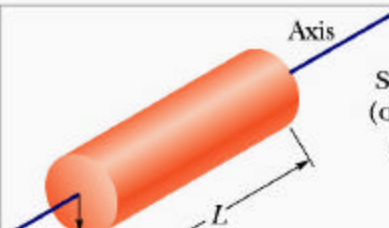
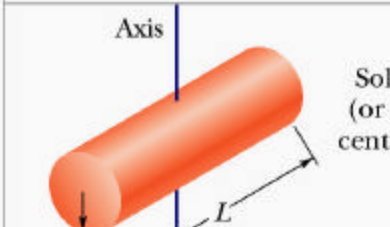
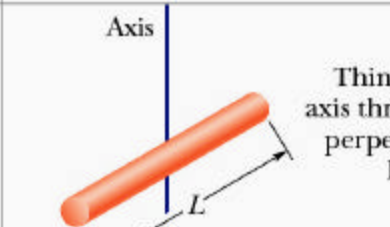
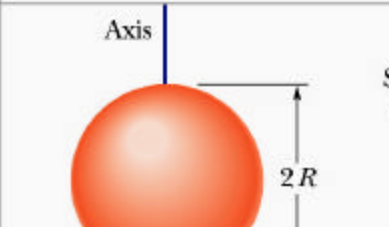
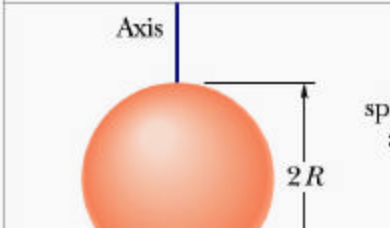
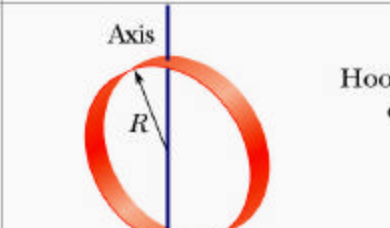
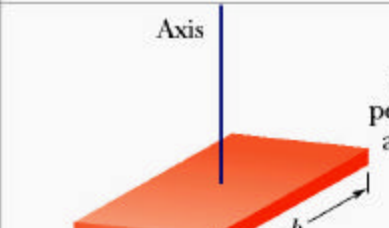
- Rod of mass,  $M$ , radius,  $R$ , and length,  $L$   $R \ll L$
- Moments of inertia about two axes, (a) and (b)

$$I_a = \frac{1}{2} MR^2$$

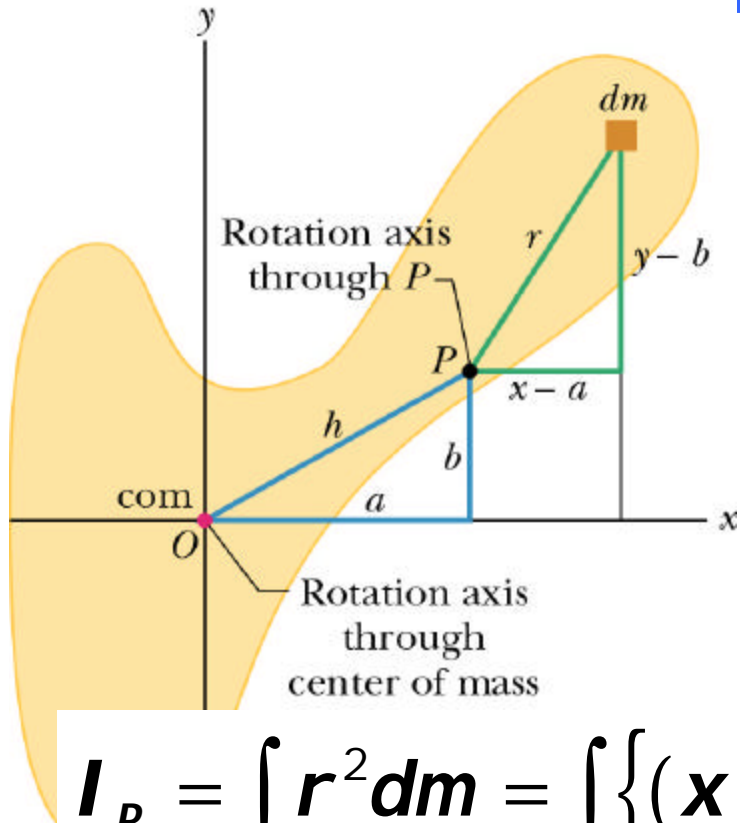
$$I_b = \frac{1}{12} ML^2$$

$$I \equiv \int_M r^2 dm$$

# Typical Shapes - table 11-2

 <p>Hoop about central axis</p> <p><math>I = MR^2</math> (a)</p>	 <p>Annular cylinder (or ring) about central axis</p> <p><math>I = \frac{1}{2}M(R_1^2 + R_2^2)</math> (b)</p>	 <p>Solid cylinder (or disk) about central axis</p> <p><math>I = \frac{1}{2}MR^2</math> (c)</p>
 <p>Solid cylinder (or disk) about central diameter</p> <p><math>I = \frac{1}{4}MR^2 + \frac{1}{12}ML^2</math> (d)</p>	 <p>Thin rod about axis through center perpendicular to length</p> <p><math>I = \frac{1}{12}ML^2</math> (e)</p>	 <p>Solid sphere about any diameter</p> <p><math>I = \frac{2}{5}MR^2</math> (f)</p>
 <p>Thin spherical shell about any diameter</p> <p><math>I = \frac{2}{3}MR^2</math> (g)</p>	 <p>Hoop about any diameter</p> <p><math>I = \frac{1}{2}MR^2</math> (h)</p>	 <p>Slab about perpendicular axis through center</p> <p><math>I = \frac{1}{12}M(a^2 + b^2)</math> (i)</p>

# Parallel Axis Theorem



- Simple technique for finding  $I$  about a different axis (when know it about axis through the cntr of mss)
- Pick origin at c of m with z-axis along w
- So  $x_{com} = y_{com} = 0$

$$I_P = \int r^2 dm = \int \left\{ (x - a)^2 + (y - b)^2 \right\} dm$$

$$= \int (x^2 + y^2) dm - 2a \int x dm - 2b \int y dm + \int h^2 dm$$

$$I_P = I_{com} + Mh^2$$



parallel axis theorem.MOV

# Torque: Can change Rotational Kinetic Energy

## Definition of Torque

$$\vec{\tau} = \vec{r} \times \vec{F}$$

$t = rF \sin \phi$  in direction  $\perp$   $rF$  plane

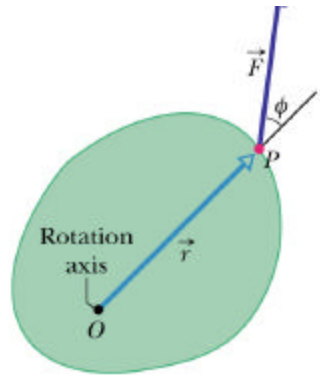
$$t = r_{\perp} F = rF_t$$

- Torque produces Angular Acceleration
- Body's rotational energy changes as result
- Work done by torque

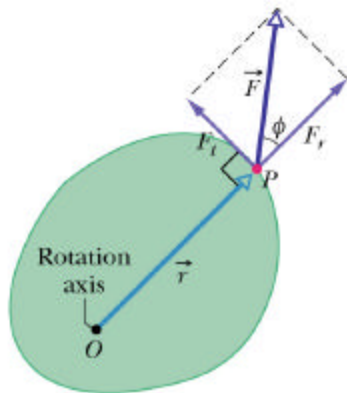
$$Work = F_t s = Frq$$

$$= tq$$

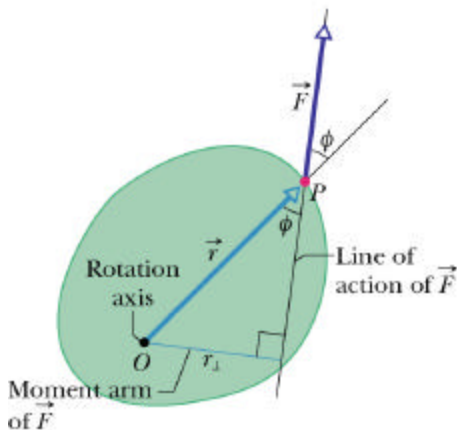
$$Power = t \frac{dq}{dt} = tw$$



(a)



(b)

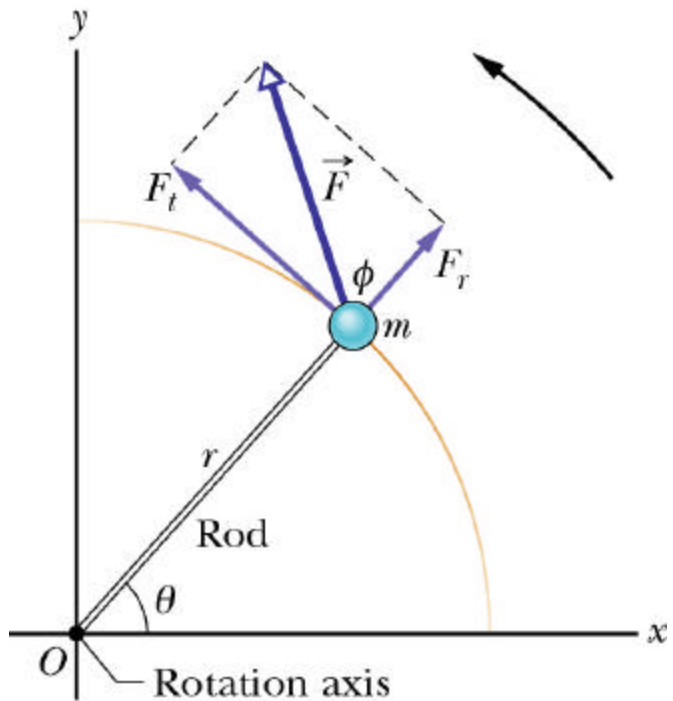


(c)



work from torque.MOV

# Utility of Torque- rotation about fixed axis



$F_r$  component can't move  $m$

$$F_t = ma_t = mra$$

$$t = rF_t = rma$$

$$t = I a \text{ with } I = mr^2$$

Note directions of  $\vec{t}$  and  $\vec{a}$

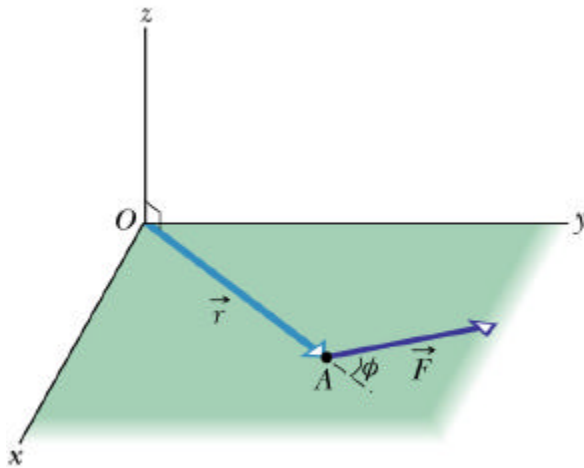
$$\vec{t} = I \vec{a} \text{ with } I = \sum m_i r_i^2$$

- Massless rod of length,  $r$ , with pt-mass,  $m$ , at end
- apply force,  $F$ , as shown

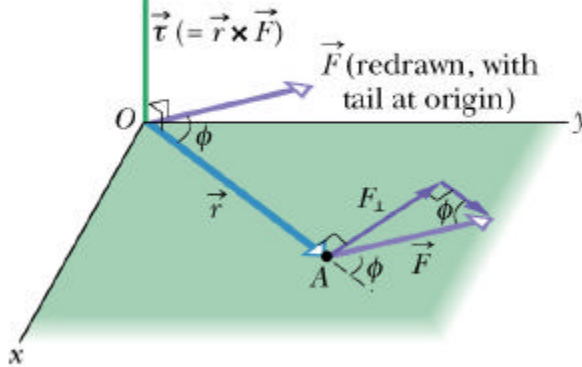
- When net torque on body is zero, there can be no angular acceleration
  - If  $w=0$ , and no torque,  $w$  stays that way

# Specifying Torque

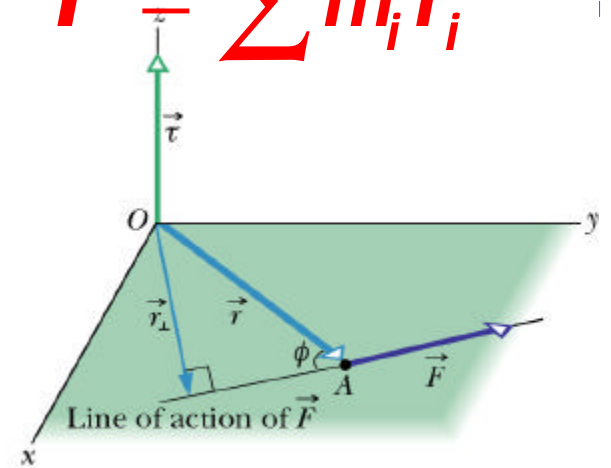
$$\vec{\tau} = I \vec{a} \quad \text{with} \quad I = \sum m_i r_i^2$$



(a)



(b)



(c)

$$\vec{\tau} = \vec{r} \times \vec{F} \quad t = rF \sin f$$

- Note the equivalence of the two perspectives

- ◆ distance times transverse force component
- ◆ lever, or moment, arm times force
- ◆ Importance of all seen in pix



platform- torque.MOV

# $I$ larger, $a$ smaller

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$$\vec{\tau} = I \vec{a} \quad \text{with} \quad I = \sum m_i r_i^2$$

- Note that this form is similar to Newton's 2<sup>nd</sup> Law:  $F=ma$ 
  - ◆ Motion driver (force) produces linear acceleration ... smaller if inertia (mass) larger
- Torque drives angular acceleration
  - ◆  $I$ , property of the body, is a measure of rotational inertia



platform- I w bodies1-SNDRS.MOV



platform- I w bodies2-SNDRS.MOV

# Torques do Work and Make Energy

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- Rotational Mechanical Energy
- Kinetic Energy result of angular acceleration or Work done
- Energy rot pix
- Analogy between energies
  - ◆ translational
  - ◆ rotational

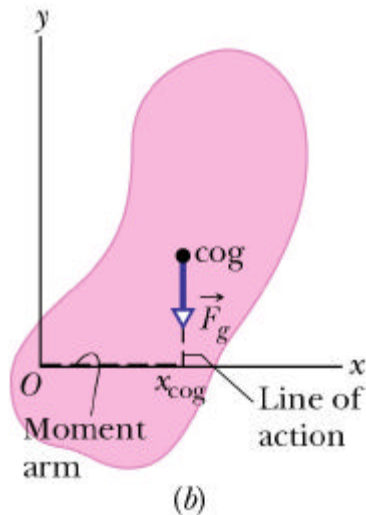
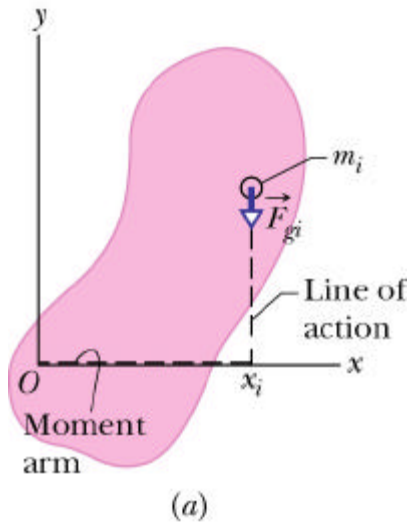
$$\begin{aligned} \textit{Work} &= F_t s = Frq \\ &= tq \end{aligned}$$

$$\textit{Power} = tw$$

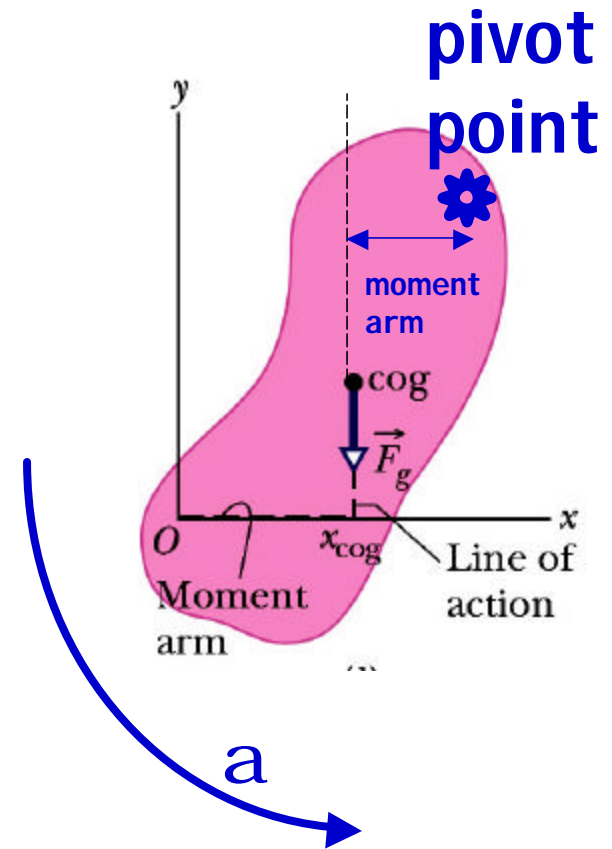


energy-power in rotational motion2.MOV

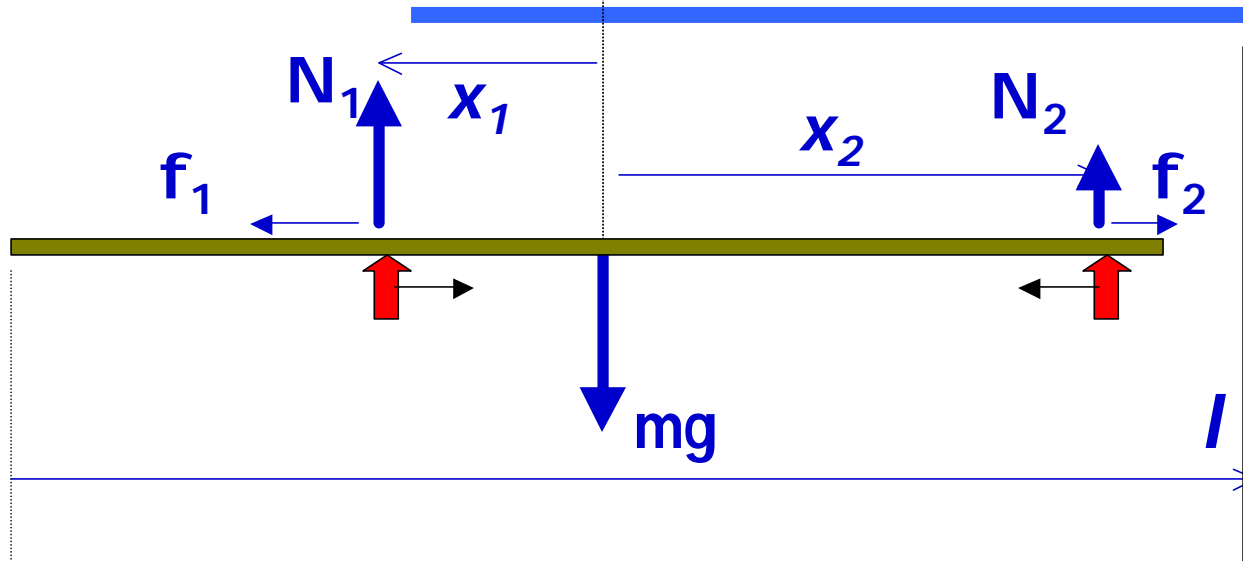
# Center of Gravity



- Rigid body feels gravity as if all the body mass concentrated at center of mass
  - ◆ See proof in text (sec 13-3)
- Consequence
  - ◆ Pivot rigid body and cog must lie below it
- Demo
- Pivot body at cog and stable (no torque)



# Finding C of M Empirically



- Allow rod to slide across hands
- Works because the forces and torques on the rod must sum to zero

Why does this work??

$$\vec{r}_{com} = \frac{1}{M} \int \vec{r} dm$$

1D -- like stick

$$m \equiv \frac{dm}{dl} = \frac{M}{L}$$

$$l_{com} = \frac{1}{M} \int_0^L l m dl = \frac{1}{L} \int_0^L l dl = \frac{L}{2}$$

# Conclude

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- Remember that HW 6 on chapters 10-11 due on Wednesday
- Read through chapters 12-13 ... will finish in the next two lectures