

Announcements

- The sample final exam and associated solutions are now posted. Check “What’s new?” at website!
- Sessions will be held to review homeworks and other course material and to go over the sample exam on Tuesday and Thursday at 3PM in room 428 Pupin.
- Special office hours will be held by me in my office on Friday from 10AM until 12:30PM (or later if people are waiting). I am likely also to be in my office next Monday morning (on the final date). I will also be available to answer any specific questions by e-mail.

Thermodynamics and Gases

Last Time

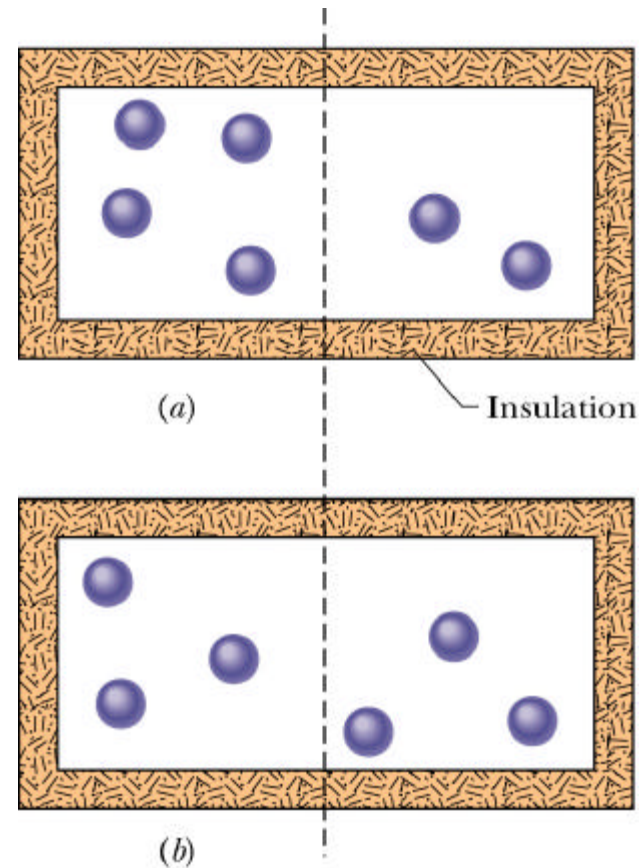
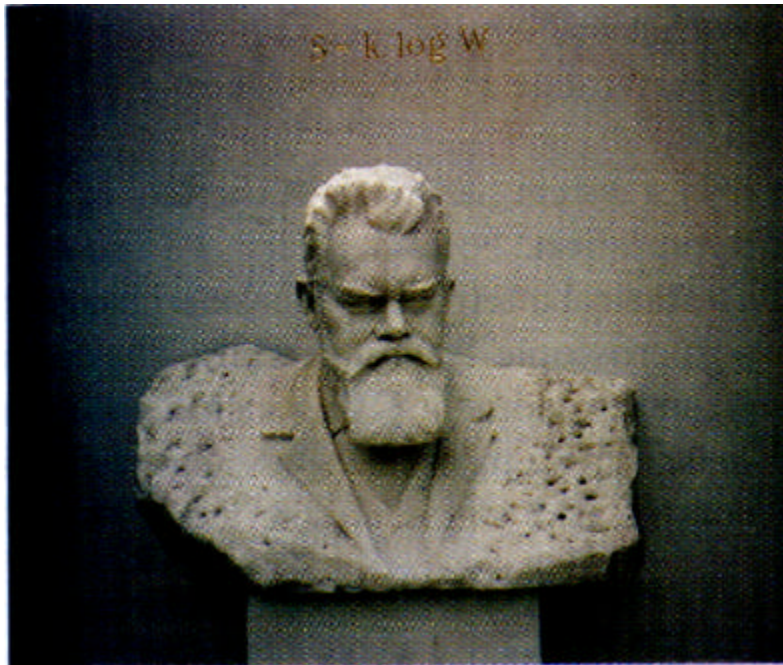
- 2nd Law of Thermodynamics
- Entropy
- Uses in engines, heat pumps, refrigerators

Today (not on final)

- Origin of Entropy
 - ⌚ measure of system disorder
 - ⌚ statistical mechanics
- ⌚ Finish

review

Boltzmann: Macrostates (and Microstates)



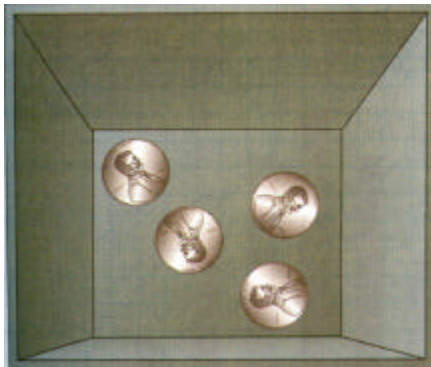
- Two scenarios at right
 - ◆ each has several possible microstates
 - ◆ depend on which molecules are on which side
- Analyze with coins

$$S = k \ln W$$

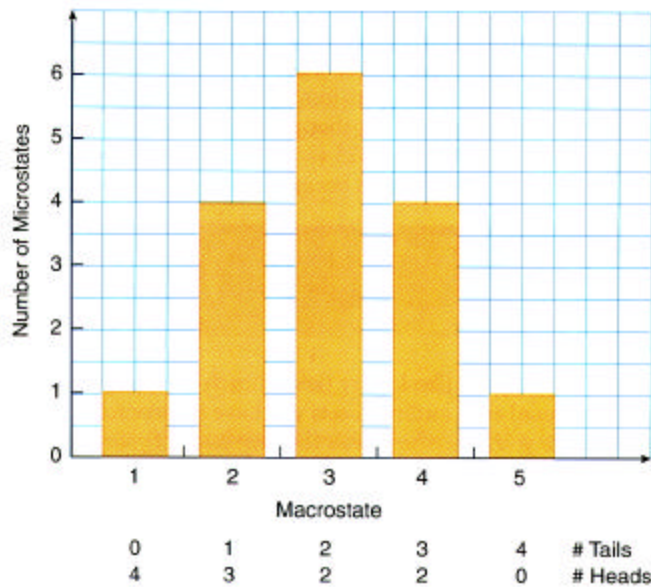
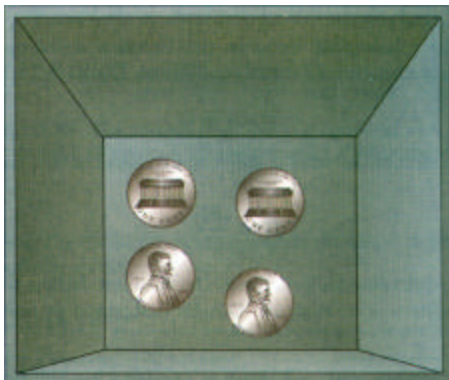
$$W = \binom{\text{number of}}{\text{microstates}}$$

review

4 pennies - macrostates and microstates



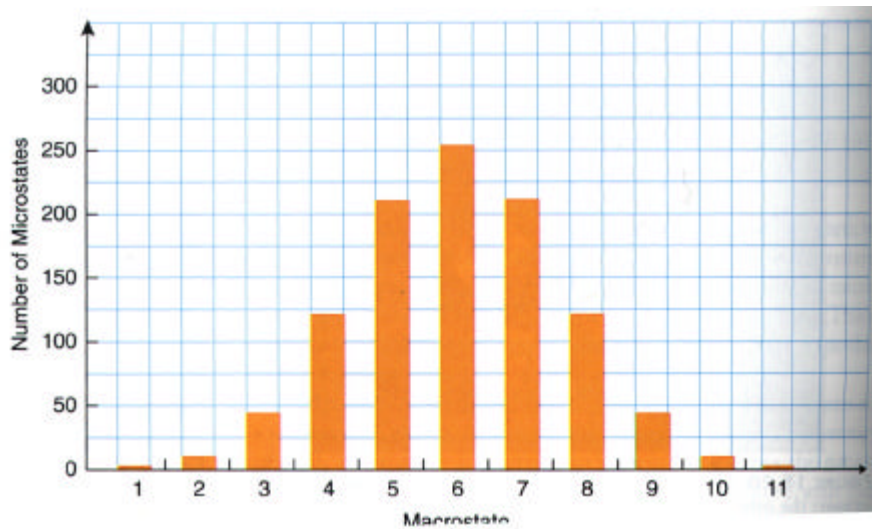
- all heads: one microstate
- One tail: 4 microstates
- 2 heads, 2 tails: six microstates
- In total: 16 microstates
- Pix shows relative probabilities



Four-Penny Macrostates and Microstates				
<i>Macrostate (1): 4 heads (H) and 0 tails (T)</i>				
<i>Penny 1</i>	<i>Penny 2</i>	<i>Penny 3</i>	<i>Penny 4</i>	
H	H	H	H	
(1 microstate)				
<i>Macrostate (2): 3 heads and 1 tail</i>				
<i>Penny 1</i>	<i>Penny 2</i>	<i>Penny 3</i>	<i>Penny 4</i>	
H	H	H	T	
H	H	T	H	
H	T	H	H	
T	H	H	H	
(4 microstates)				
<i>Macrostate (3): 2 heads and 2 tails</i>				
<i>Penny 1</i>	<i>Penny 2</i>	<i>Penny 3</i>	<i>Penny 4</i>	
H	H	T	T	
H	T	T	H	
T	T	H	H	
T	H	H	T	
T	H	T	H	
H	T	H	T	
(6 microstates)				
<i>Macrostate (4): 1 head and 3 tails</i>				
<i>Penny 1</i>	<i>Penny 2</i>	<i>Penny 3</i>	<i>Penny 4</i>	
H	T	T	T	
T	H	T	T	
T	T	H	T	
T	T	T	H	
(4 microstates)				
<i>Macrostate (5): 0 heads and 4 tails</i>				
<i>Penny 1</i>	<i>Penny 2</i>	<i>Penny 3</i>	<i>Penny 4</i>	
T	T	T	T	
(1 microstate)				

review

10 coins - contrast



Macrostate	H	T	# microstates Ω
1	10	0	1
2	9	1	10
3	8	2	45
4	7	3	120
5	6	4	210
6	5	5	252
7	4	6	210
8	3	7	120
9	2	8	45
10	1	9	10
11	0	10	1
Total			$1024 = 2^{10}$

- 5 each is 252 times more probable than all heads
- as more coins are added, equal partitioning = equal numbers heads and tails gets more probable relative to all other possibilities
 - ◆ especially those near the extremes (list top and bottom)

Increasing the coins

TABLE 15-2
Probabilities of Various
Macrostates for 100 Coin
Tosses

Macrostate		Number of Microstates, W
Heads	Tails	
100	0	1
99	1	1.0×10^2
90	10	1.7×10^{13}
80	20	5.4×10^{20}
60	40	1.4×10^{28}
55	45	6.1×10^{28}
50	50	1.0×10^{29}
45	55	6.1×10^{28}
40	60	1.4×10^{28}
20	80	5.4×10^{20}
10	90	1.7×10^{13}
1	99	1.0×10^2
0	100	1

- As the number of coins increases, equal numbers as heads/tails is more and more probable relative to others
 Note: binomial dist.

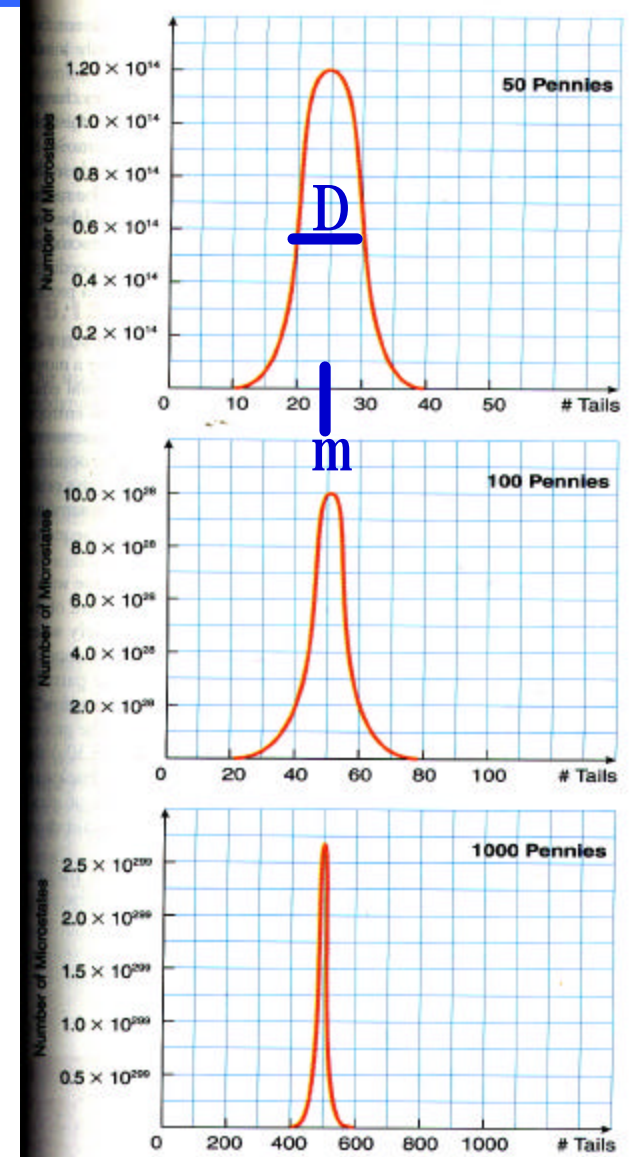
$$W = \frac{N!}{n_1! n_2!}$$

Ratio width to mean
 of parameter

$$\frac{\Delta}{m} \propto \frac{1}{\sqrt{N}}$$

With atoms we deal
 with $\sim 10^{24}$ atoms

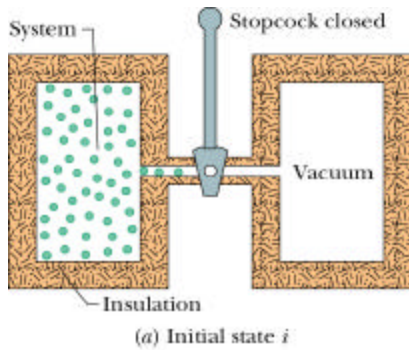
Frank Sciulli



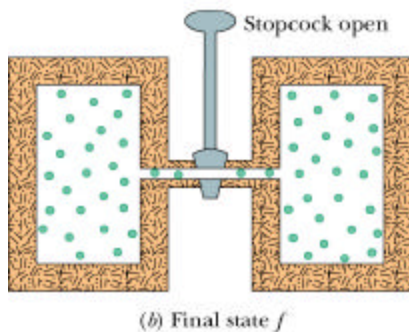
slide 6

review

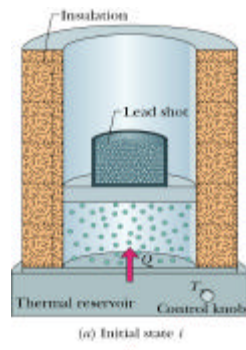
Entropy increase of free expansion of gas from Thermodynamics



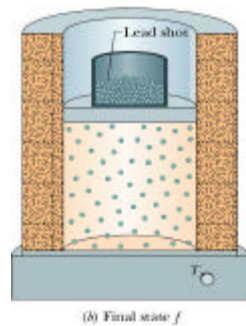
Irreversible process



irreversible
in a closed
system



Reversible process



reversible
isothermal
equivalent

$$DS = \frac{Q}{T} = \frac{W}{T} = \frac{nRT}{T} \ln \frac{V_f}{V_i}$$

$$DS = nR \ln \frac{V_f}{V_i}$$

- This is the entropy change of the gas for both processes

more general

entropy change for ideal gas

$$DS = \int_i^f \frac{dQ}{T} = nC_V \int_{T_i}^{T_f} \frac{dT}{T} + nR \int_{V_i}^{V_f} \frac{dV}{V}$$

$$DS = nC_V \ln \frac{T_f}{T_i} + nR \ln \frac{V_f}{V_i}$$

Boltzmann applied = Statistical Mechanics

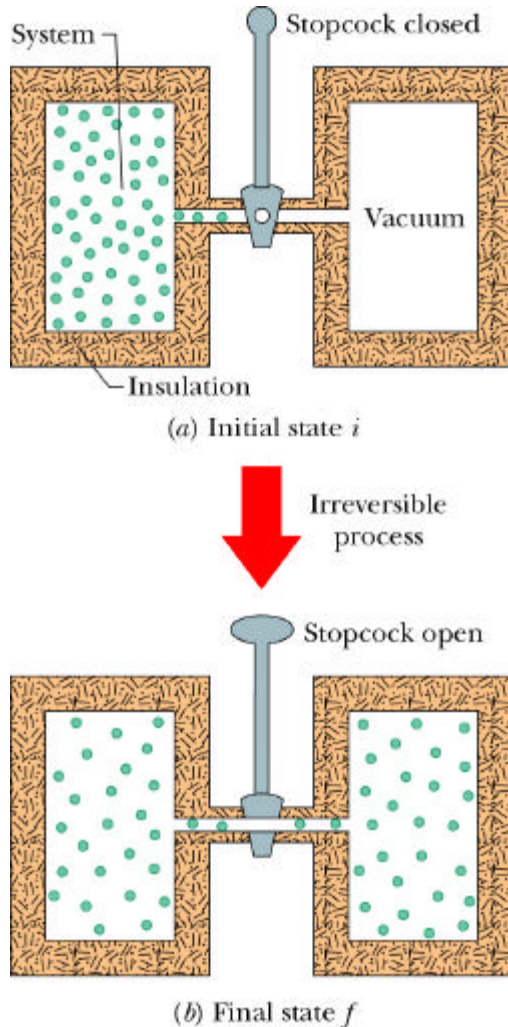
- More microstates = more entropy
- For atoms and microstates, the number of states are delineated by
 - ◆ Position possibilities
 - ◆ Velocity possibilities
- for each atom, enormous number of possibilities
- Total entropy = Sum over all atoms (also enormous number)
- Have looked at simpler physical system: limited number of coins - states of up or down

$$S_{molecule} = k \ln w$$

w = # microscopic states in macrostate

$$S_{gas} = \sum_{\text{all molecules}} S_{molecule} \quad \Delta S_{1 \rightarrow 2} = Nk [\ln w_2 - \ln w_1] = Nk \ln \left(\frac{w_2}{w_1} \right)$$

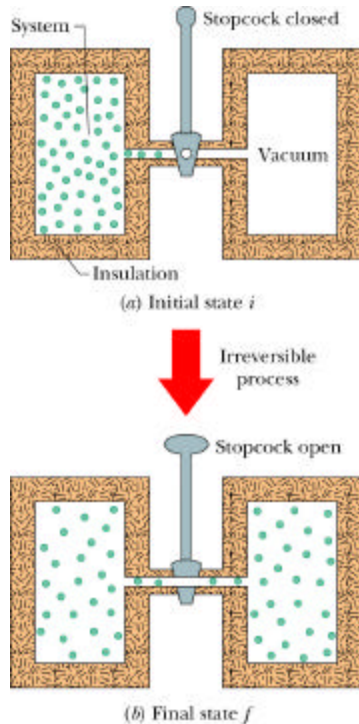
Position states available depends on Volume



- Pix shows gas expanding to new volume: getting many more position states available
- In general, same atom in volume V_f has many more position states available than at volume V_i . ($w \propto V$)
- Ratio of number of available position states is V_f / V_i .

$$\frac{\text{No states } V_f}{\text{No states } V_i} = \frac{\int_{V_i}^{V_f} \frac{1}{V} dV}{\int_{V_i}^{V_i} \frac{1}{V} dV} = \ln \frac{V_f}{V_i}$$

Entropy of Isothermal Expansion from Stat Mech



- Ratio of number of position states is V_f / V_i .

$$\frac{\text{No states } V_f}{\text{No states } V_i} = \frac{\omega(V_f)}{\omega(V_i)}$$

$$DS_{j^{\text{th}} \text{ molecule}} = k \ln \frac{\omega(V_f)}{\omega(V_i)}$$

$$DS = \sum_{j=1}^N DS_j = Nk \ln \frac{\omega(V_f)}{\omega(V_i)}$$

$$DS = nR \ln \frac{\omega(V_f)}{\omega(V_i)}$$

entropy change for
isothermal gas expansion
from using thermodynamics

$$DS = nR \ln \frac{V_f}{V_i}$$

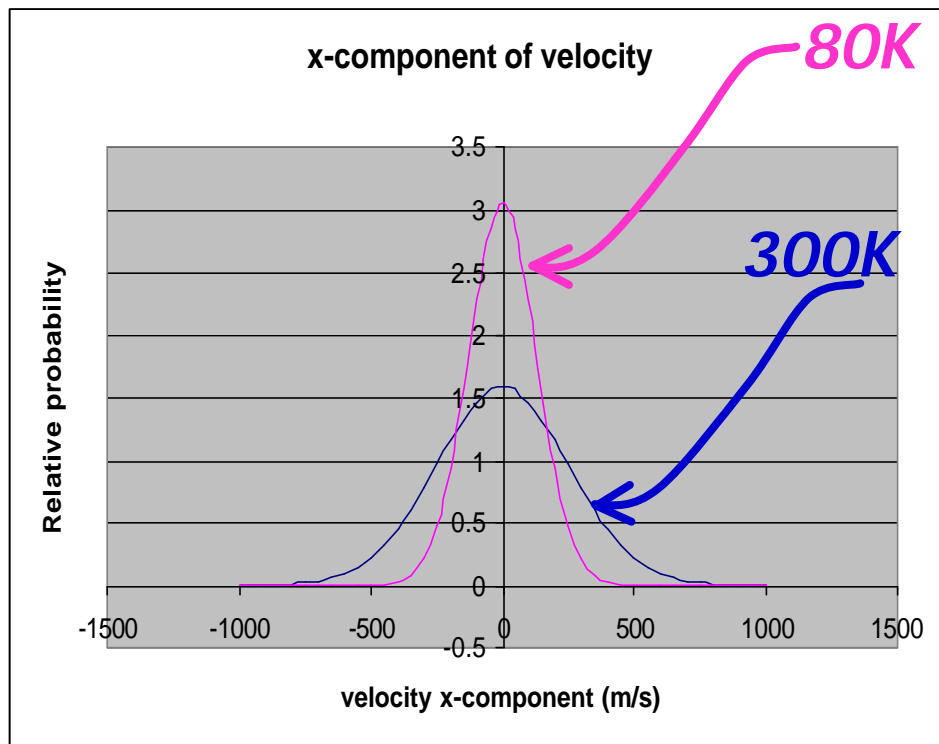
Temp Changes: Velocity States Available depends on Temperature

$$dP(v_x) = \frac{1}{\sqrt{2\pi kT}} e^{-\frac{mv_x^2}{2kT}} dv_x \quad \text{Gaussian}$$

$$\langle v_x \rangle = 0$$

$$S_{v_x} = \sqrt{\frac{kT}{m}}$$

- Number of v_x states available to noble gas atom is proportional to S_{v_x}
- Number of available velocity states proportional to $S_{v_x} * S_{v_y} * S_{v_z}$



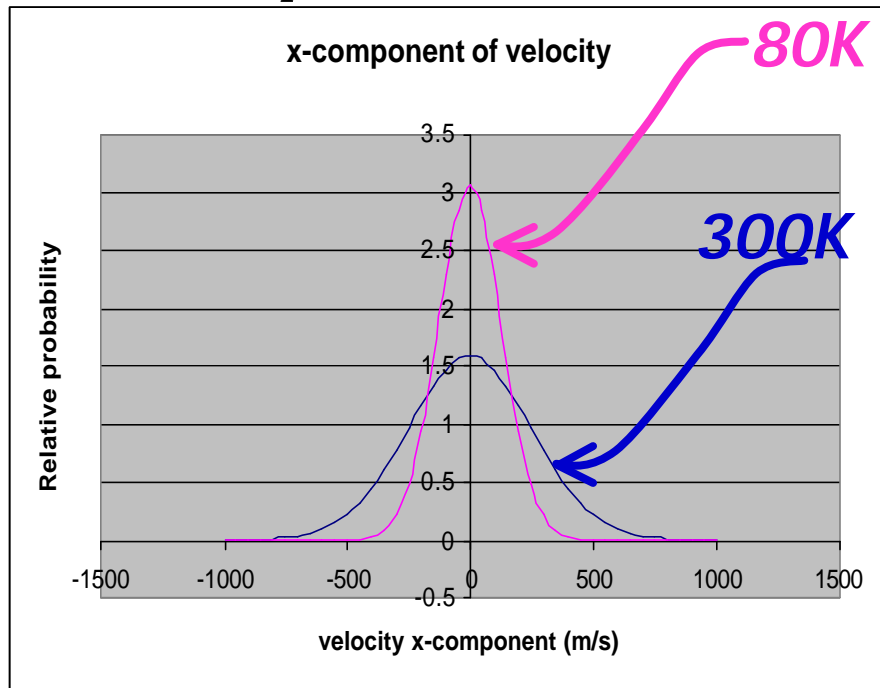
$$S_{v_x} = \sqrt{\frac{kT}{m}}$$

$$S_{v_x} S_{v_y} S_{v_z} = \left(\frac{kT}{m}\right)^{3/2}$$

$$\frac{\text{No states } T_f}{\text{No states } T_i} = \left(\frac{T_f}{T_i}\right)^{3/2}$$

Entropy Change for Temperature Change of monatomic gas from Stat Mech

$$dP(v_x) = \frac{1}{\sqrt{2\pi kT}} e^{-\frac{mv_x^2}{2kT}} dv_x \quad \text{Gaussian}$$



entropy change for isochoric change from using thermodynamics

$$DS = nC_V \ln \frac{T_f}{T_i}$$

- Number of available velocity states proportional to $S_{v_x} * S_{v_y} * S_{v_z}$

$$S_{v_x} S_{v_y} S_{v_z} = \left(\frac{2\pi kT}{m} \right)^{3/2}$$

$$\frac{\text{No states } T_f}{\text{No states } T_i} = \left(\frac{T_f}{T_i} \right)^{3/2}$$

$$DS_{1 \text{ molecule}} = k \ln \left(\frac{T_f}{T_i} \right)^{3/2}$$

$$DS = Nk \ln \left(\frac{T_f}{T_i} \right)^{3/2}$$

$$DS = \frac{3}{2} Nk \ln \left(\frac{T_f}{T_i} \right) = nC_V \ln \left(\frac{T_f}{T_i} \right)$$

Check Stat Mech vs Thermo

Stat mech-Vol. chg

$$\frac{\text{No states } V_f}{\text{No states } V_i} = \frac{\alpha V_f^3}{\alpha V_i^3}$$

$$DS_{1 \text{ molecule}} = k \ln \frac{\alpha V_f^3}{\alpha V_i^3}$$

$$DS = Nk \ln \frac{\alpha V_f^3}{\alpha V_i^3}$$

$$DS = nR \ln \frac{\alpha V_f^3}{\alpha V_i^3}$$

Stat mech = temp chg

$$S_{V_x} S_{V_y} S_{V_z} = \left(\frac{kT}{m} \right)^{3/2}$$

$$\frac{\text{No states } T_f}{\text{No states } T_i} = \left(\frac{\alpha T_f^3}{\alpha T_i^3} \right)^{3/2}$$

$$DS_{1 \text{ molecule}} = k \ln \left(\frac{\alpha T_f^3}{\alpha T_i^3} \right)^{3/2}$$

$$DS = Nk \ln \left(\frac{\alpha T_f^3}{\alpha T_i^3} \right)^{3/2}$$

$$DS = \frac{3}{2} Nk \ln \frac{\alpha T_f^3}{\alpha T_i^3}$$

$$DS = nC_V \ln \frac{\alpha T_f^3}{\alpha T_i^3}$$

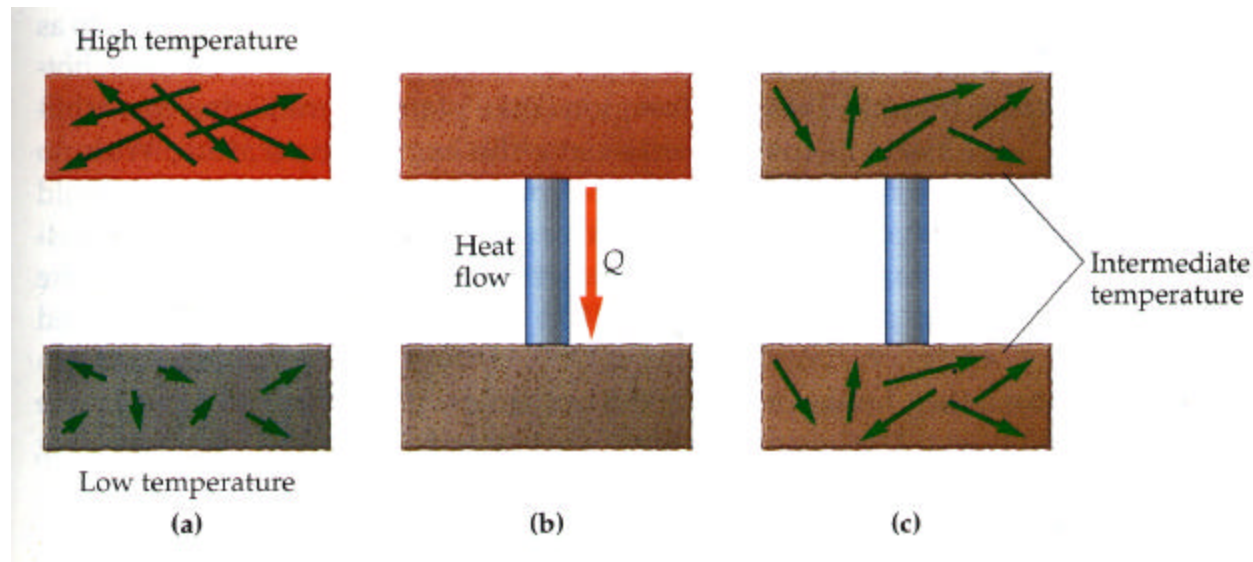
thermodynamics
entropy change for ideal gas

$$DS = \int_i^f \frac{dQ}{T} = nC_V \int_{T_i}^{T_f} \frac{dT}{T} + nR \int_{V_i}^{V_f} \frac{dV}{V}$$

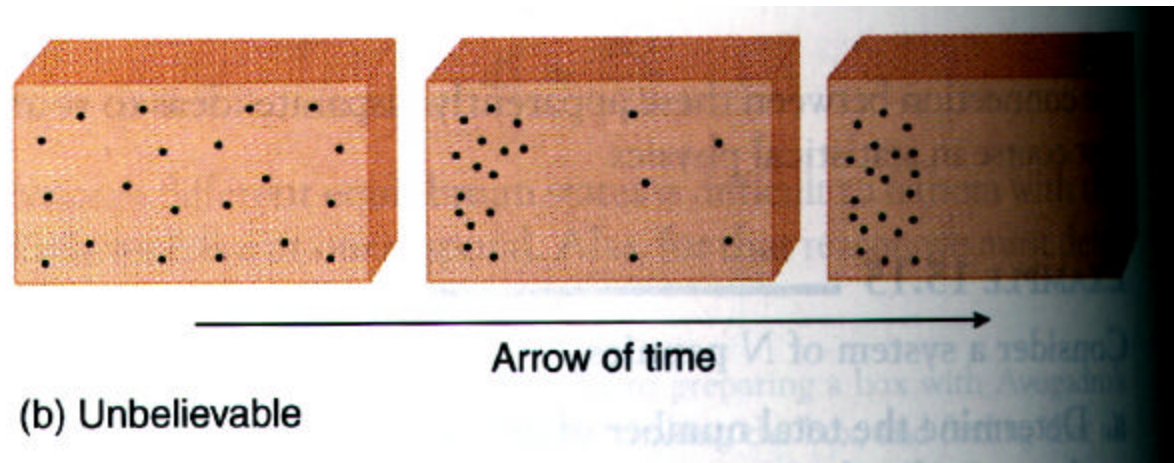
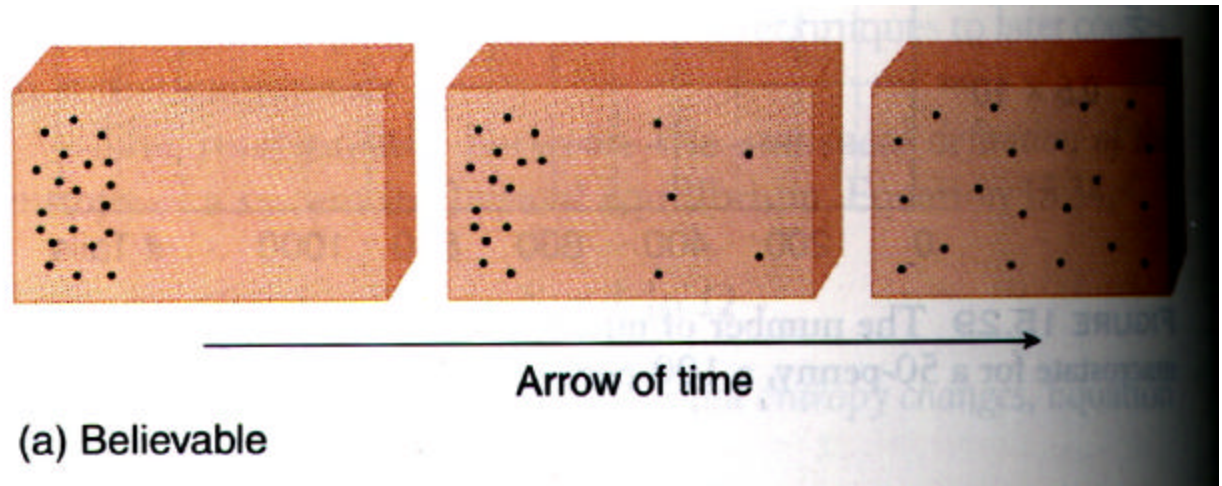
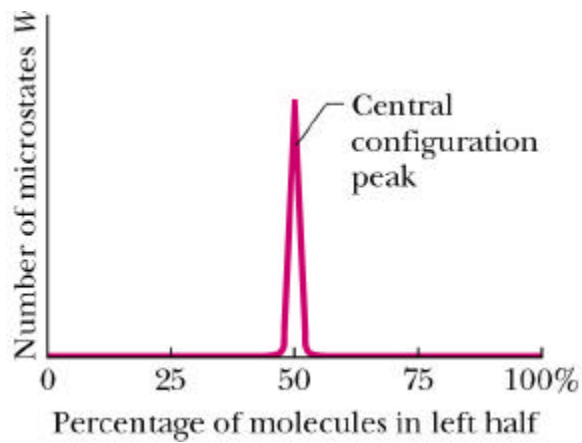
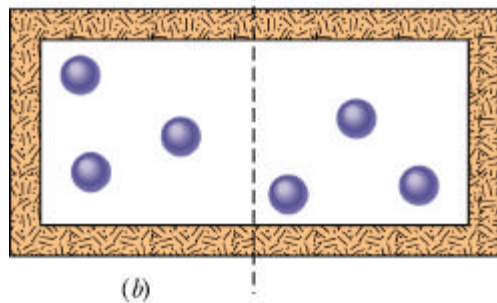
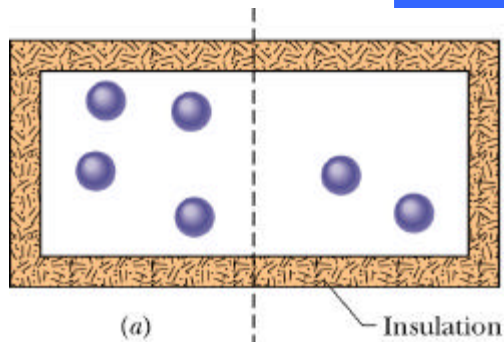
$$DS = nC_V \ln \frac{T_f}{T_i} + nR \ln \frac{V_f}{V_i}$$

Statistical Requirement

- Systems will evolve from less probable macrostate to most probable macrostate
- Entropy will always increase
- Ordered segregation will inexorably become disordered equilibrium

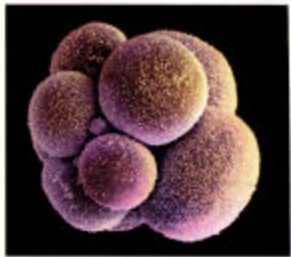


Molecules in a box



Exceptions?

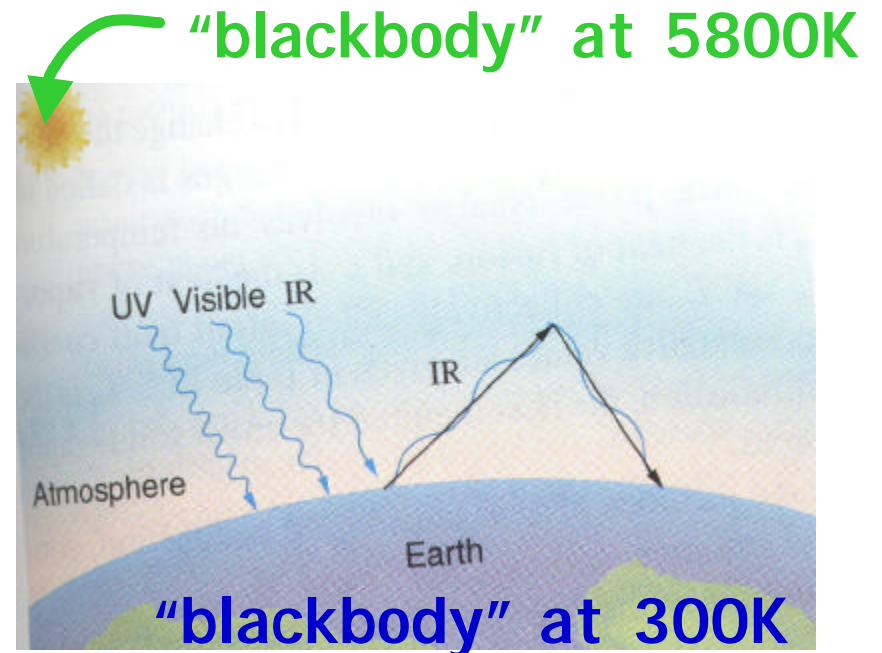
- Do living organisms violate entropic principles?
- Does evolution of life violate entropic principles?



NO!

Entropy of the world increases in every process!

Evolution of life does not contradict physical principles!



That's Almost All

- Last lecture
- Different lecturer next semester
- Before leave ... one more story of personal nature (with a moral)

1 - wildlife



2 - bird feeder & birds



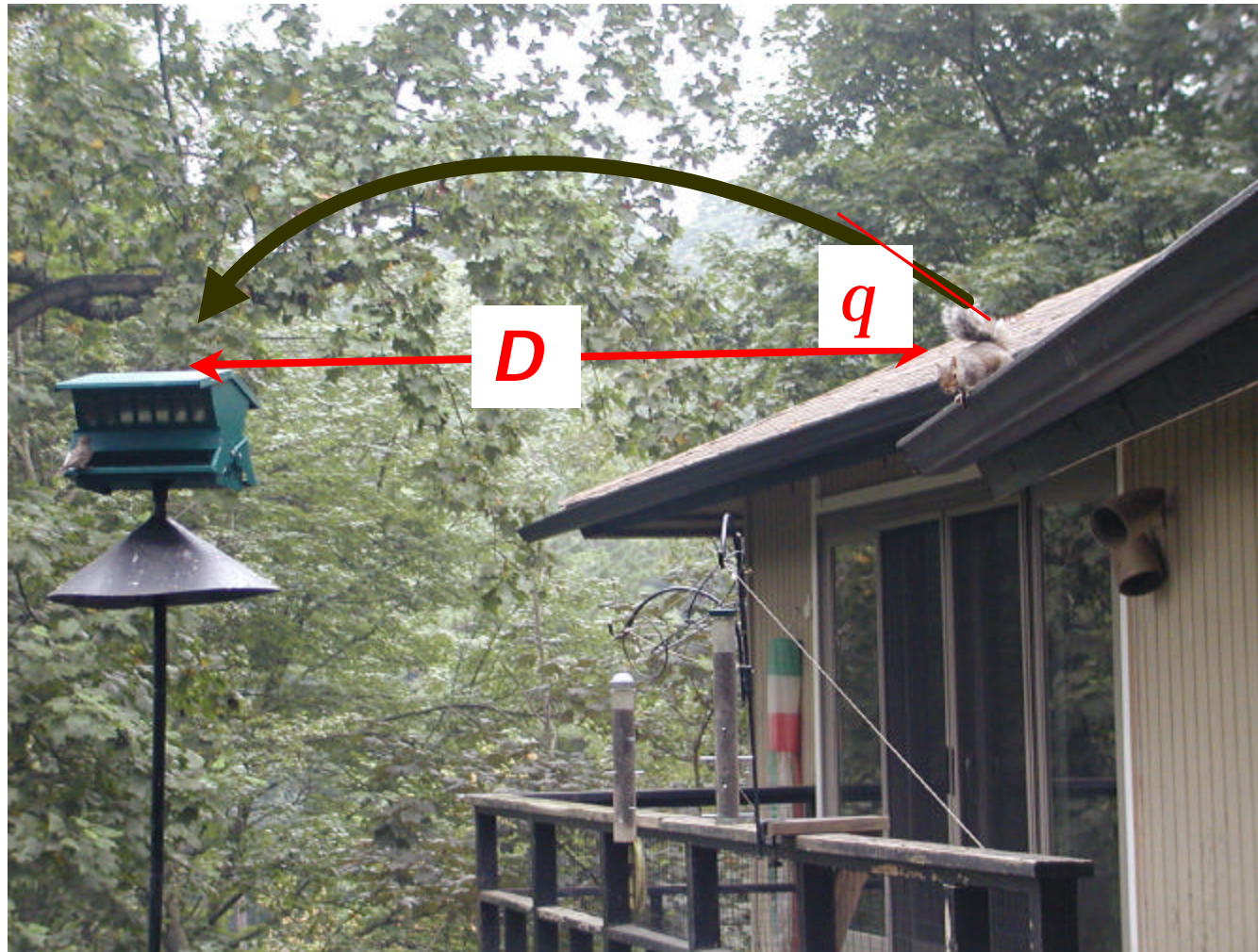
3 - who's in charge?



4 - not me, that's for sure!



5 - how does squirrel make it?



Implications for the jump:
In simple model, the coefficient of kinetic friction at the feeder must be large enough to stop the horizontal momentum of the squirrel

$$m_k > \frac{1}{\tan q}$$

Solution: make top of feeder slippery

6 - They still get there!



7 - An experimental result



Moral of the Story

- Learn all the established principles
- Know where they come from and how to apply them
- **But check them out**
 - ◆ **With experiments**
 - ◆ **With observations**



Surrender... Solution: no seed in the feeder!

- no food
- no squirrels
- no birds

- unhappy birds
- unhappy squirrels

9: Revenge

- tap, tap, tap
- on my chamber door
- outside my window



That's All

- **Have a good final!**
- **Have a good second semester!**
- **Best of fortune in the future!**
- **Please remain to fill in Questionnaire**